



Lecture 7

Pumping & Population Inversion*

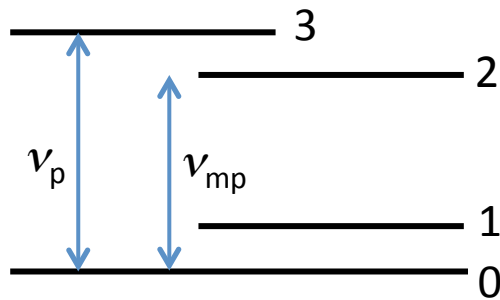
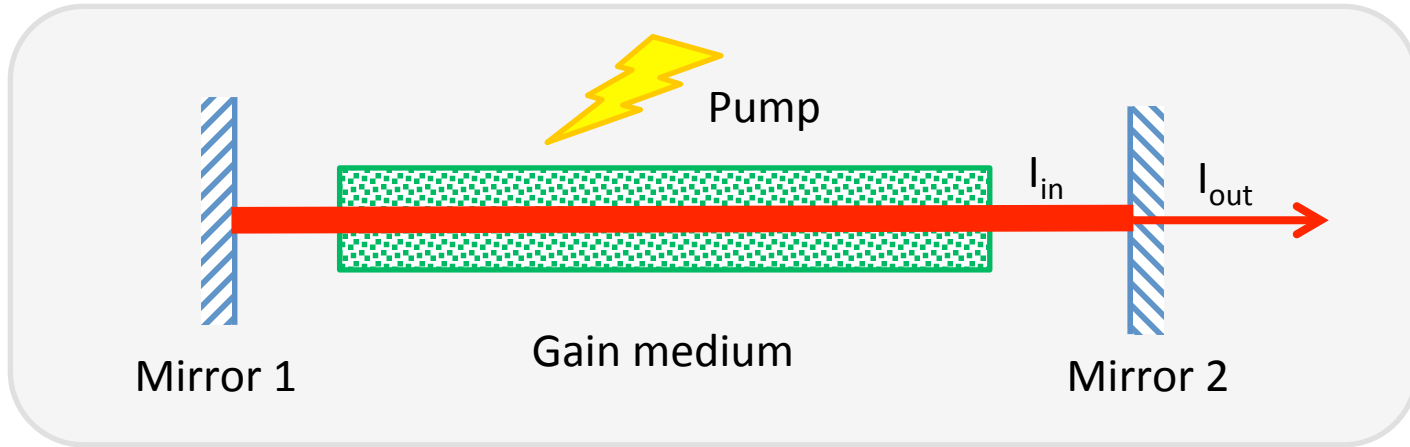
Min Yan

Optics and Photonics, KTH

Reading

- *Principles of Lasers* (5th Ed.): Chapter 6.
- Skip: Quantum mechanical treatment in 6.4.1.1
- Squeeze: 6.3.3-6.3.5, 6.4.1-6.4.4.

Laser



- Pump efficiency: $\eta_p = P_m / P_p$
- Pump rate: R_p , or dN_2/dt
- Threshold pump power: P_{th}

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4. Electrical pumping	25'
Total:	80'

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Pumping varieties

- **Optical pumping** (CW/pulsed)
 - **Lamp:** for solid-state/liquid lasers [broad absorption bands]
 - **Laser:** for solid-state/liquid/gas lasers [broad/narrow absorption bands/lines]
- **Electrical pumping** (CW/RF/pulsed; for gas/semiconductor lasers)
- **X-ray pumping**
- **E-beam pumping**
- **Chemical pumping**

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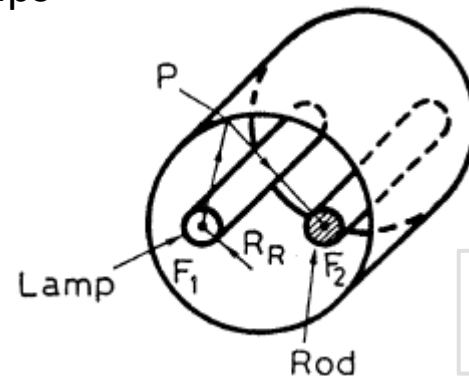
Pumping: Lamp

Lamp types

- *For pulsed lasers:*
Medium-to-high pressure (500~1500 Torr) Xe or Kr flashlamps
- *For CW lasers:*
High-pressure (1-8atm) Kr lamps

- **Elliptical cylinder**

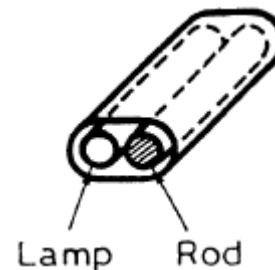
- Rod and lamp at foci
- Specular reflection



Rod radius: mm~cm
Rod length: cm~<1m

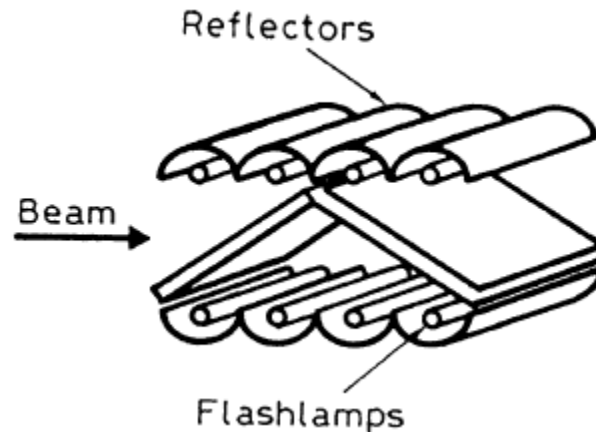
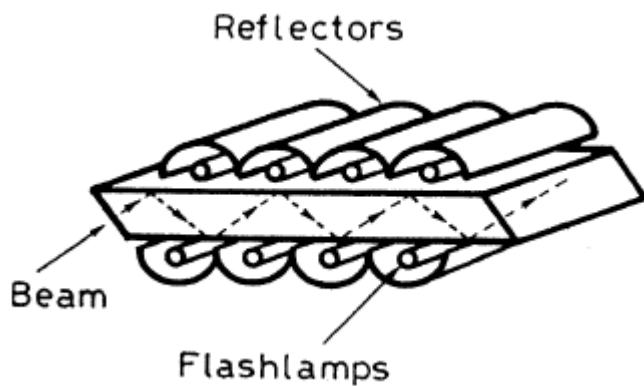
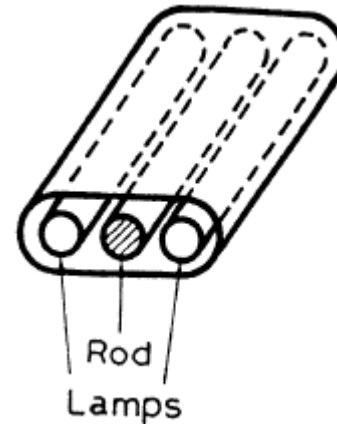
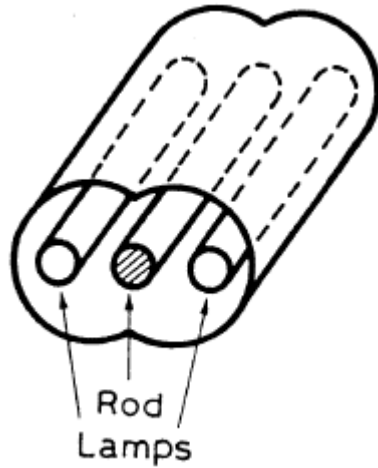
- **Close-coupling**

- Close-pack
- Diffusive reflection
- ✓ More uniform



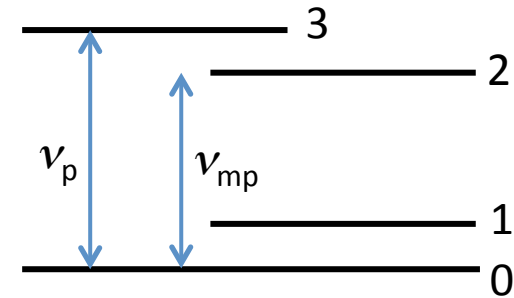
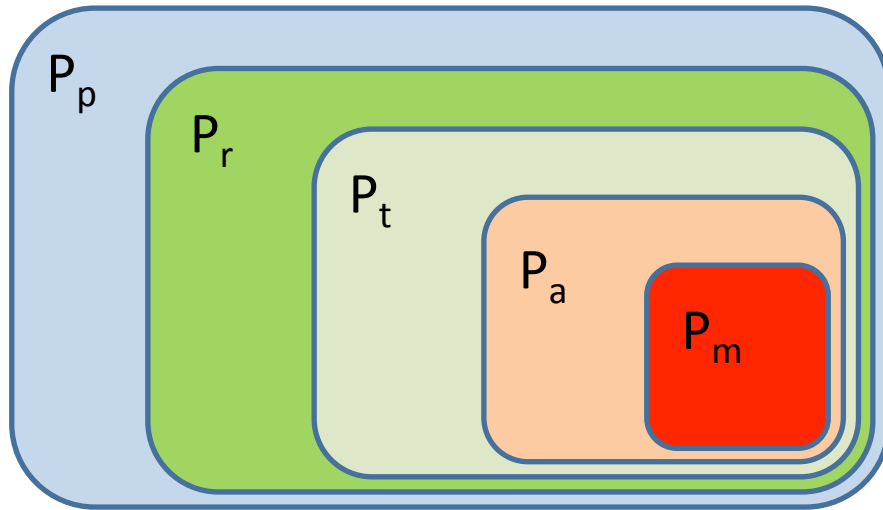
Pumping: Lamps

- $P_p \uparrow \eta_p \downarrow$
- More uniform



Pump efficiency and pump rate

Assumption: uniform R_p

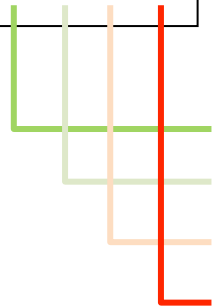


$$P_m = \left(\frac{dN_2}{dt} \right)_p V h \nu_{mp} = R_p V h \nu_{mp}$$

$$P_m = P_p \eta_p$$

$$R_p = \eta_p \frac{P_p}{A h \nu_{mp}}$$

$$\eta_p = \frac{P_r}{P_p} \frac{P_t}{P_r} \frac{P_a}{P_t} \frac{P_m}{P_a} = \eta_r \eta_t \eta_a \eta_{pq}$$



- Radiative efficiency
- Transfer efficiency
- Absorption efficiency
- Power quantum efficiency

Pump efficiency, η_p

TABLE 6.1. Comparison between computed pumping efficiency terms for different laser materials

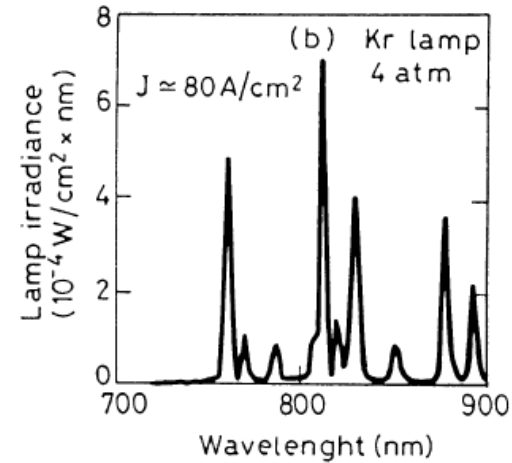
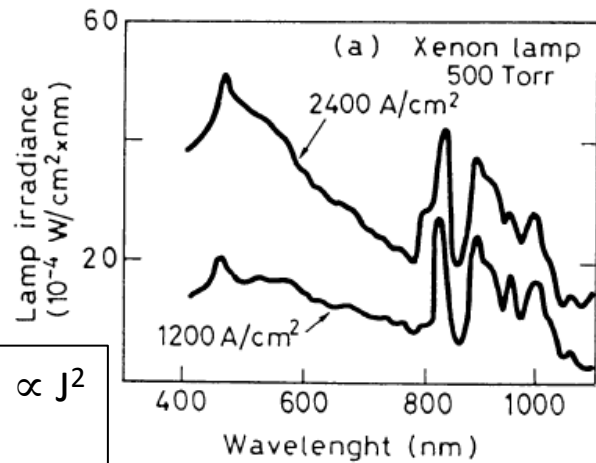
Active Medium	η_r (%)	η_t (%)	η_a (%)	η_{pq} (%)	η_p (%)
Ruby	27	78	31	46	3.0
Alexandrite	36	65	52	66	8.0
Nd:YAG	43	82	17	59	3.5
Nd:Glass (Q-88)	43	82	28	59	5.8

Configuration: Elliptical cylinder

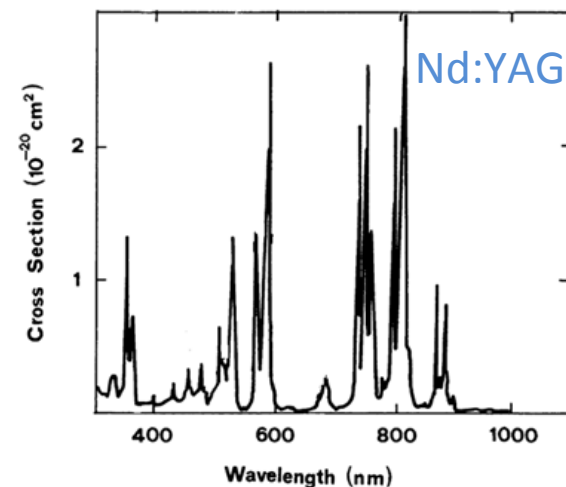
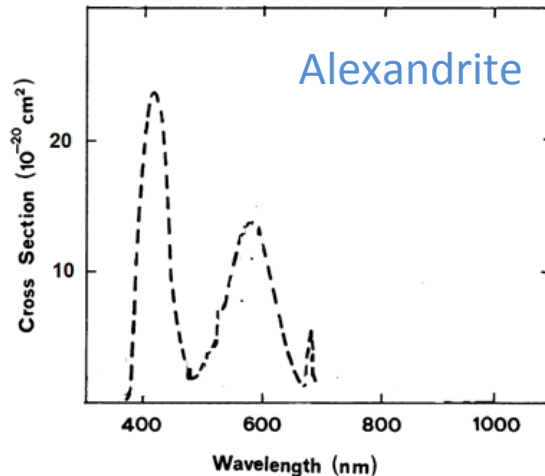
Absorption efficiency, η_a

Emission

- Continuum $\propto J^2$
- Lines $\propto J$



Absorption



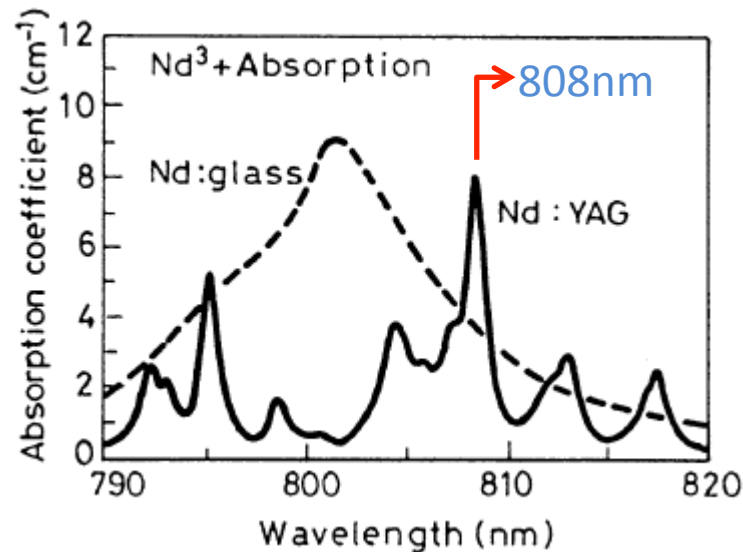
- Nd:YAG: Nd³⁺ in Y₃Al₅O₁₂ crystal (host-independent)
- Alexandrite: Cr³⁺ in BeAl₂O₄ crystal (host-dependent)

Contents

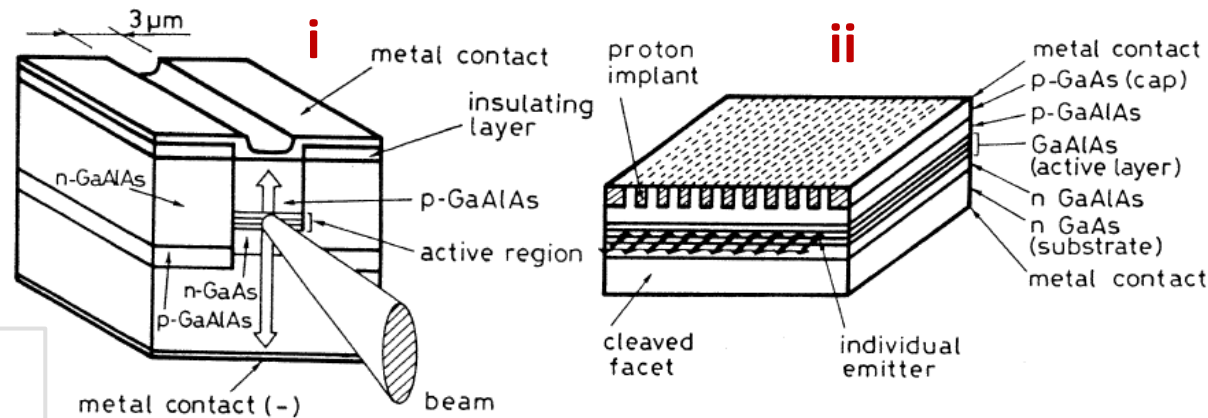
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Laser pumping

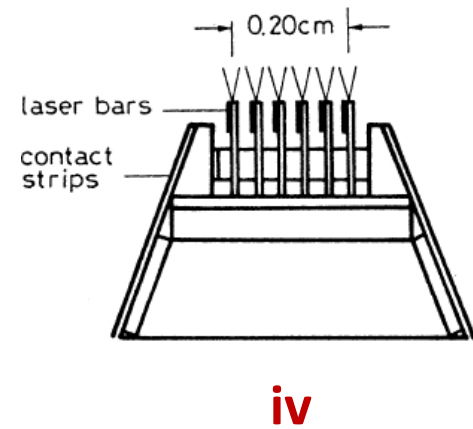
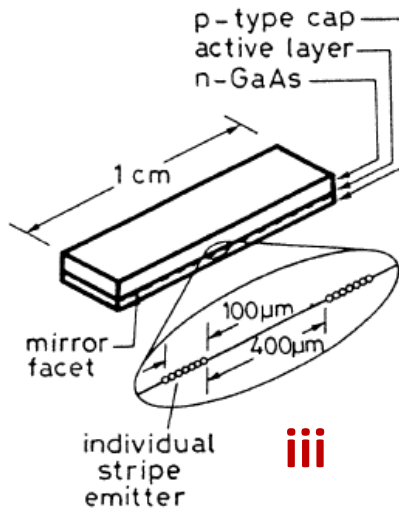
- **Convenience:** Efficient and high-power diode lasers widely available
- **Key examples:**
 - **Nd:YAG** and “siblings” pumped by GaAs/AlGaAs QW-lasers (~800nm)
 - Yb:YAG, Nd:glass or Yb:Er:glass pumped by InGaAs/GaAs QW-lasers (950~980nm)
 - Alexandrite, Cr:LISAF pumped by GaInP/AlGaInP QW-lasers (640-680nm)
 - Tm:Ho:YAG laser pumped by AlGaAs QW-lasers (785nm)



Pumping diodes

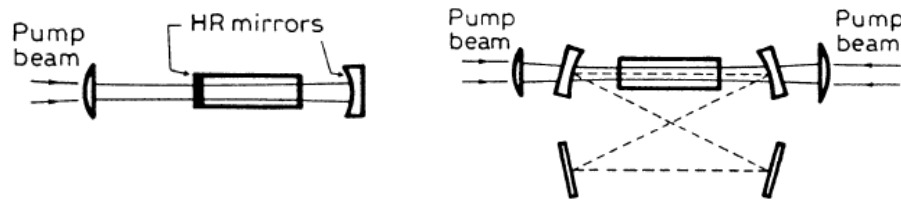


- i. Single-diode
Beam: $2 \times 4 \mu\text{m}^2$; $P < 100 \text{mW}$
- ii. Diode-array
 $P \approx 2 \text{W}$
- iii. Diode-bar
 $P = 10\text{-}20 \text{W}$
- iv. Stacked-bars
 $P = 100 \text{W}$; emission bandwidth \uparrow

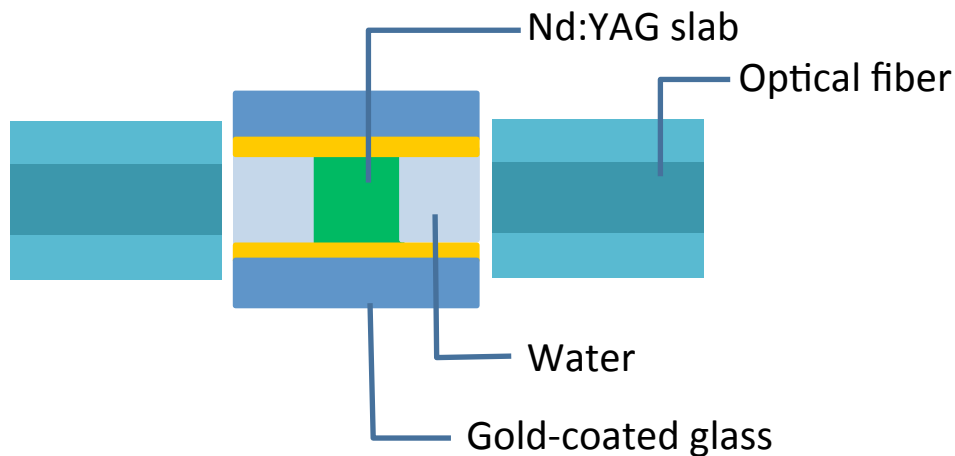


Pumping scheme

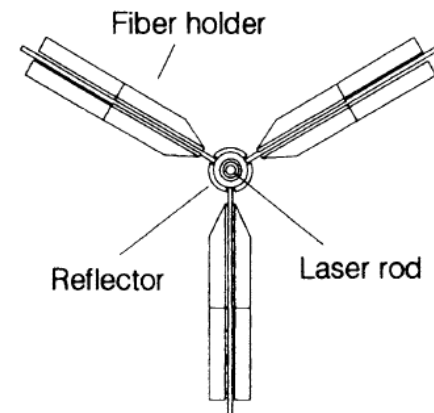
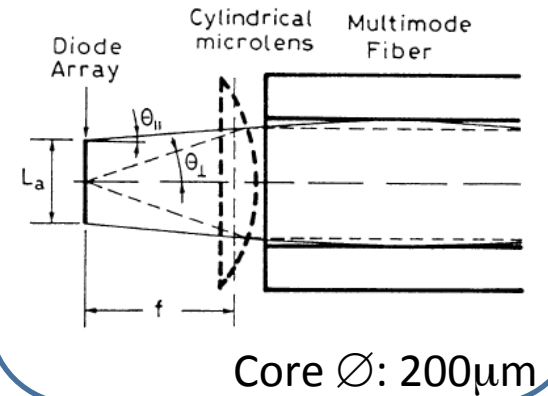
- Longitudinal pumping



- Transverse pumping



Beam shaping



10W CW

Pump efficiency

Nd:YAG

TABLE 6.3. Comparison between pumping efficiencies of lamp pumping and diode pumping

Pump Configuration	η_r (%)	η_t (%)	η_a (%)	η_{pq} (%)	η_p (%)
Lamp	43	82	17	59	3.5
Diode (longitudinal)	50	80	98	82	32
Diode (transverse)	50	80	90	82	30

- **Lamp v.s. Diode:** η_a (6 \times) and η_{pq} (1.5 \times)

Pump rate, $\langle R_p \rangle$

For uniform R_p

$$R_p = \eta_p \frac{P_p}{A l h \nu_{mp}}$$

Space-dependent lasing/pump beams: a matter of calculating A

1. **Diode (longitudinal)** $\langle R_P \rangle = \eta_{pd} \frac{P_p}{h \nu_{mp}} \frac{2}{\pi (w_0^2 + w_p^2) l}$

η_{pd} : η_p for diode pumping

Optimum: $w_p \approx w_0$

2. **Diode (transverse)** $\langle R_P \rangle = \eta_{pd} \frac{P_p}{h \nu_{mp}} \frac{1 - \exp\left(-\frac{2a^2}{w_0^2}\right)}{\pi a^2 l}$

Optimum a exists

3. **Lamp** $\langle R_P \rangle = \eta_{pl} \frac{P_p}{h \nu_{mp}} \frac{1 - \exp\left(-\frac{2a^2}{w_0^2}\right)}{\pi a^2 l}$

η_{pl} : η_p for lamp pumping

w_0 : Laser beam waist
 w_p : pump beam waist
 a : radius of the active region

Assumption for cases 2 and 3:

- $R_p = \text{const}$ if $r < a$ (doped region)
- $R_p = 0$ if $r > a$ (undoped cladding)

Threshold pump power, P_{th}

Procedure:

$$\langle N_2 \rangle_c \rightarrow \langle R_p \rangle_c \rightarrow P_{th}$$

$$\langle N_2 \rangle_c = \frac{\gamma}{\sigma_e l}$$

$$\langle R_p \rangle_c = \frac{\langle N_2 \rangle_c}{\tau}$$

γ : logarithmic cavity loss

σ_e : effective stimulated emission cross-section

l : active medium length

τ : upper-level lifetime

Comment: $\langle N_2 \rangle_c$ differs for quasi-3-level system

1. Diode (longitudinal)
$$P_{th} = \frac{\gamma}{\eta_{pd}} \frac{h\nu_{mp}}{\tau} \frac{\pi(w_0^2 + w_p^2)}{2\sigma_e}$$

2. Diode (transverse)
$$P_{th} = \frac{\gamma}{\eta_{pd}} \frac{h\nu_{mp}}{\tau} \frac{\pi a^2}{\sigma_e \left[1 - \exp\left(-\frac{2a^2}{w_0^2}\right) \right]}$$

- Can be 50% larger than case 1

3. Lamp
$$P_{th} = \frac{\gamma}{\eta_{pl}} \frac{h\nu_{mp}}{\tau} \frac{\pi a^2}{\sigma_e \left[1 - \exp\left(-\frac{2a^2}{w_0^2}\right) \right]}$$

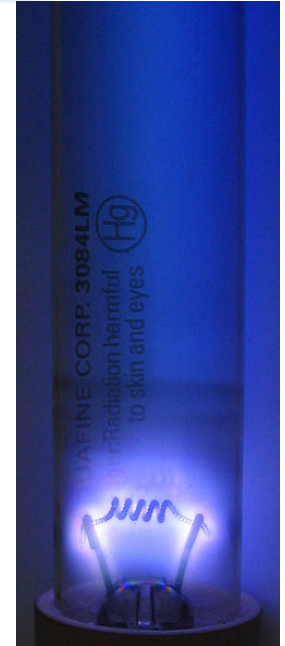
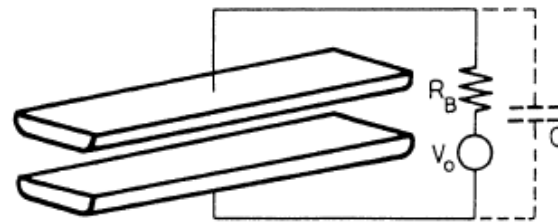
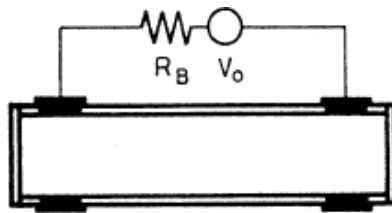
- Can be 10 times larger than case 2
- High thermal load

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Electrical pumping

- **Applicability:** Gas or semiconductor lasers
- **Principle:** Gas discharge
- **Configuration:**
 - *Longitudinal:* Uniform pumping (glow discharge)
 - *Transverse:* small voltage (arc discharge @ edges)



Low-pressure mercury vapor discharge [Wikipedia]

He

Ne

Ar

Kr

Xe

N

O₂

H₂



Excitation mechanisms

- **Collision of the first kind:**



- Or: **electron-impact excitation**
- For single-species gas
- More common
- Non-resonant
- Prelude for 2nd-kind collision

Excited state:

- Vibrational
- Rotational
- Electronic

- **Collision of the second kind:**



- Resonant
- $\Delta E < kT$ such that probability is high

Pump rate

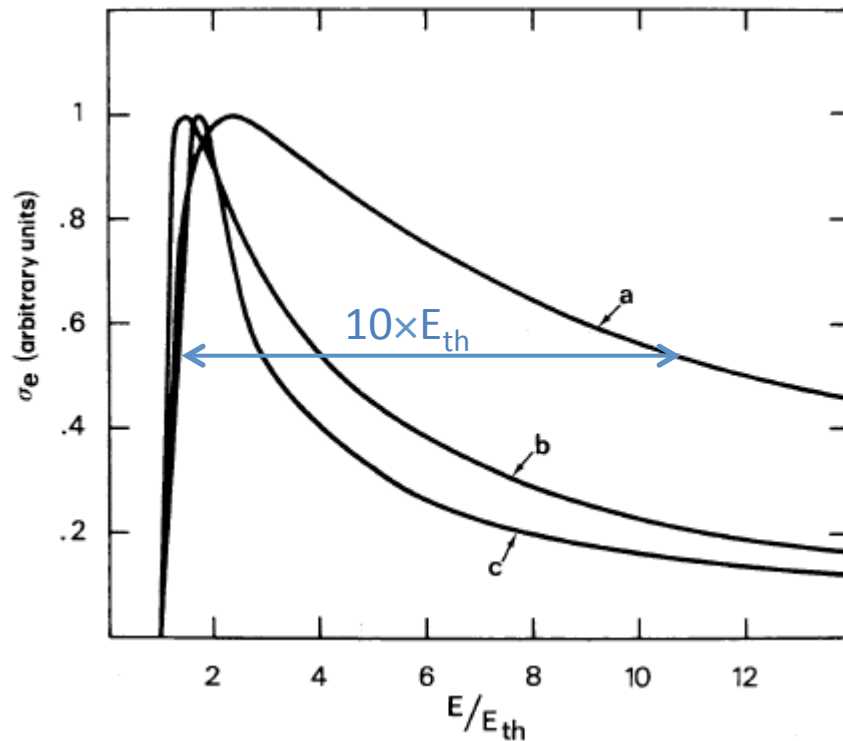
$$R_p = \left(\frac{dN_2}{dt} \right)_p = \sigma_{e2} N_t N_e v$$

- σ_{e2} : Electronic excitation cross-section
 - Dependent on electron energy E, as $\sigma_{e2}(E)$
- N_t : Total species population
- N_e : Electron density
- v : Electron velocity
 - Follow a distribution function $f(E)$

$$R_p = N_t N_e \langle v \sigma_{e2} \rangle$$

σ_{e2}

- a) Optically allowed transition
- b) Optically forbidden transition involving no change of multiplicity
- c) Optically forbidden transition involving a change of multiplicity



$f(E)$: Electron energy distribution

- $\mathbf{v} = \mathbf{v}_{th} + \mathbf{v}_{drift}$

Electric-field-induced velocity

Thermal velocity: random, temperature-dependent, $\langle v \rangle$

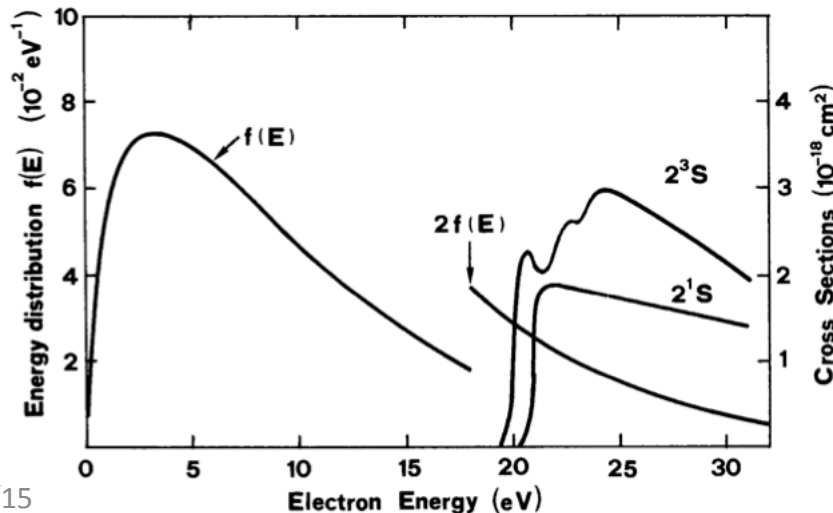
$$v_{drift}/v_{th} \approx 0.01$$

- $f(E)$: Maxwell-Boltzmann distribution

$$f(E) = \frac{2}{\sqrt{\pi}kT_e} \sqrt{\frac{E}{kT_e}} \exp\left(-\frac{E}{kT_e}\right)$$

T_e : Electron temp.

$$T_e \rightarrow f(E) \rightarrow v_{th} \rightarrow \langle v \rangle$$



He-Ne:

- $f(E)$: Maxwellian (mean electron energy of 10 eV);
- σ_{e2} : $1^1S \rightarrow 2^1S$; $1^1S \rightarrow 2^3S$ transitions of He;
- High electron energy; **inefficient**

Balance conditions

T_e is related to electrical field and pressure

$$T_e = f\left(\frac{\mathcal{E}}{p}\right)$$

- **Energy-balance condition:**

Energy loss by electron collisions = Energy supplied to electrons by E field

- **Momentum-balance condition:**

Momentum should be conserved with a collision

T_e is related to gas pressure and tube radius

$$f(T_e) = \frac{Const.}{(pR)^2}$$

- **Ionization balance condition:**

Ionization \rightleftharpoons Electron-ion recombination at tube walls

Scaling law of gas laser

$$T_e = f\left(\frac{\mathcal{E}}{p}\right)$$

$$f(T_e) = \frac{\text{Const.}}{(pR)^2}$$

Optimum T_e (for achieving max R_p) can be obtained with various combinations of R , p and \mathcal{E} .

Example: $R \downarrow \rightarrow p \uparrow \rightarrow \mathcal{E} \uparrow$

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